



# X-Ray and Neutron Science

7<sup>th</sup> ESRF/ILL International  
Student Summer Programme

5<sup>th</sup> September – 1<sup>st</sup> October 2021



# Principles of synchrotron radiation

Nicola Carmignani

ESRF, Beam Dynamics group, Accelerator and Source division

With help from Boaz Nash, Andrea Franchi, Jean-Luc Revol and Simone Liuzzo



The European Synchrotron

## **Properties of radiation**

### **Where x-rays come from**

**review of relativity**

**bending magnet radiation**

**undulator radiation**

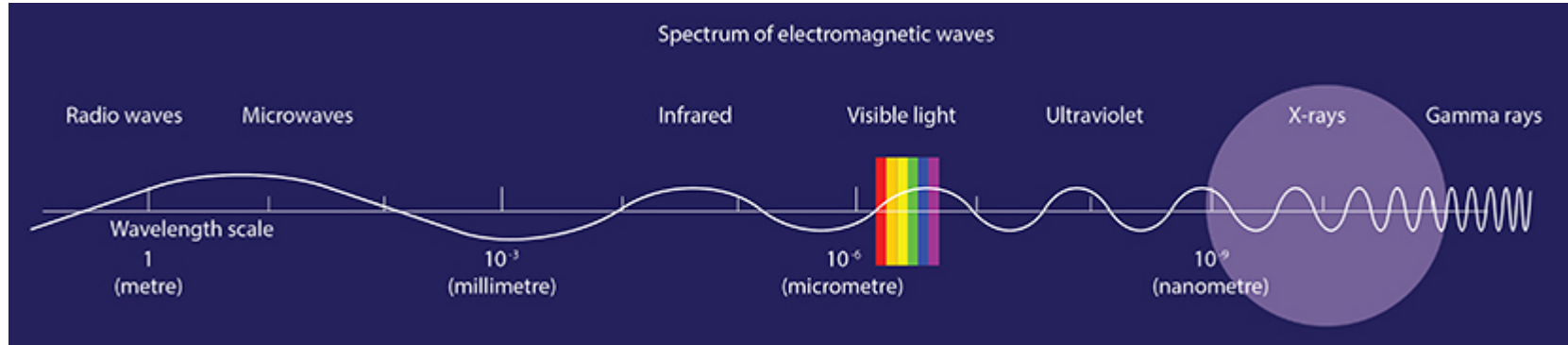
### **Particle accelerators:**

**how to store electron beams**

**what is the distribution of the electron beam**

### **ESRF upgrade: EBS**

## Spectrum of electro-magnetic radiation



$$f = c / \lambda$$

$$E_{\gamma} = h f$$

$$c = 299792458 \text{ m/s}$$

$$h = 4.135667517 \times 10^{-15} \text{ eV s}$$

$\lambda$  is the photon wavelength

$f$  is the photon frequency

$E_{\gamma}$  is the photon energy

$c$  is the speed of light

$h$  is the Planck constant

$E_{\gamma}$  for X-rays is about 1-100 keV

**Flux** of radiation:  
number of photons per second

**Brightness** of radiation:  
flux at a specific wavelength divided by  
source size and divergence

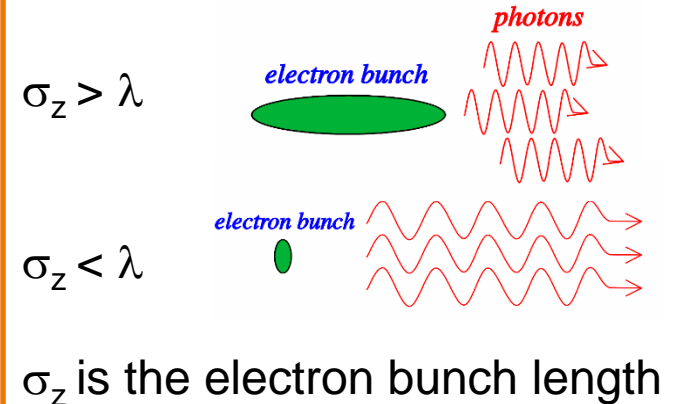


**Low brightness: low  
photon density on sample**



**High brightness: high  
photon density on sample**

**Coherence** of radiation:  
are the photons produced in  
phase or not?



**Polarization** of radiation:  
directionality of radiation field  
linear, circular partial/full  
polarization

Synchrotron light sources give some control over all these properties

Accelerating charged particles emit electro-magnetic radiation

Maxwell's equations:

$$\nabla \cdot \mathbf{E} = 4\pi\rho$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \frac{1}{c} \left( 4\pi\mathbf{J} + \frac{\partial \mathbf{E}}{\partial t} \right)$$

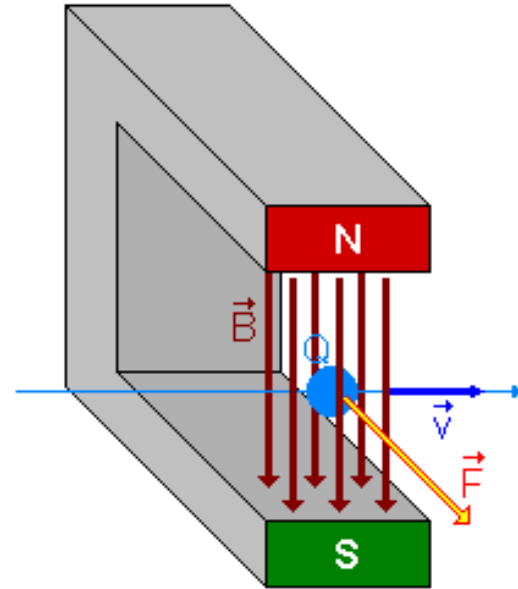
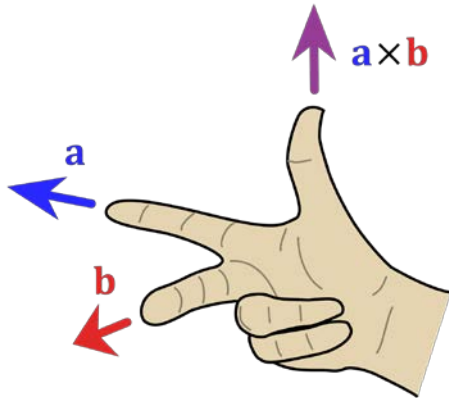
Non-relativistic Larmor formula:

$$P = \frac{q^2 a^2}{6\pi\epsilon_0 c^3}$$

P is the power radiated  
q is the electric charge  
a is the acceleration

We can move charged particles using the Lorentz force

$$\vec{F} = q \cdot (\vec{E} + \vec{v} \times \vec{B})$$



The force, and so the acceleration, is perpendicular to the particle velocity. Particles travel on an arc of circumference inside the dipole magnets.

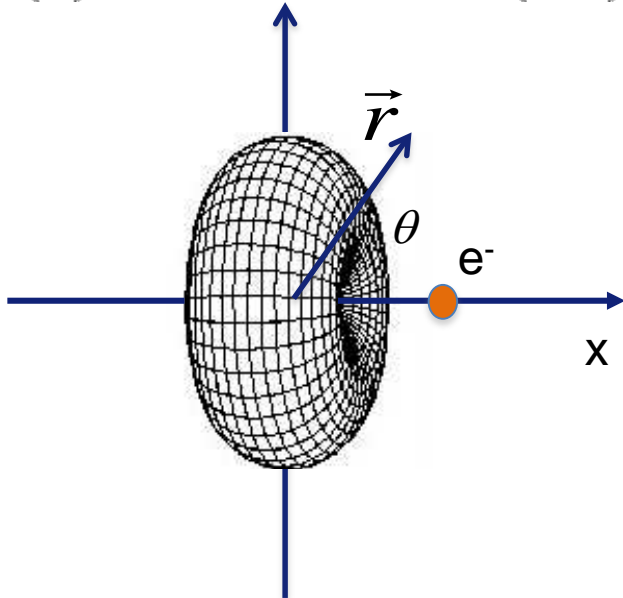
# SINGLE PARTICLE IN A HARMONIC OSCILLATOR

Charge oscillating along x

$$x(t) = A \cos(\omega t)$$

$$v(t) = -A\omega \sin(\omega t)$$

$$a(t) = -A\omega^2 \cos(\omega t)$$



It emits radiation in all directions

$$\langle S \rangle \propto \frac{\sin^2 \theta}{r^2} \hat{r}$$

S is the Poynting vector (energy per unit of time and surface)

Frequency of radiation is

$$f = \frac{\omega}{2\pi}$$

Wavelength of radiation is

$$\lambda = \frac{c}{f}$$



$$\beta = \frac{v}{c}$$

$$\gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

**c** = 299792458 m/s  
is the speed of light

$$\text{Total energy } E = \gamma mc^2$$

$$\text{Kinetic energy } E_k = (\gamma - 1)mc^2$$

$$E_k \rightarrow \frac{1}{2}mv^2, \text{ for } \beta \ll 1$$

For an electron with energy E:

$$\gamma = \frac{E}{mc^2} = \frac{E}{510998.9461 \text{ eV}}$$

$$m_{e^-} = 510998.9461 \frac{\text{eV}}{c^2}$$



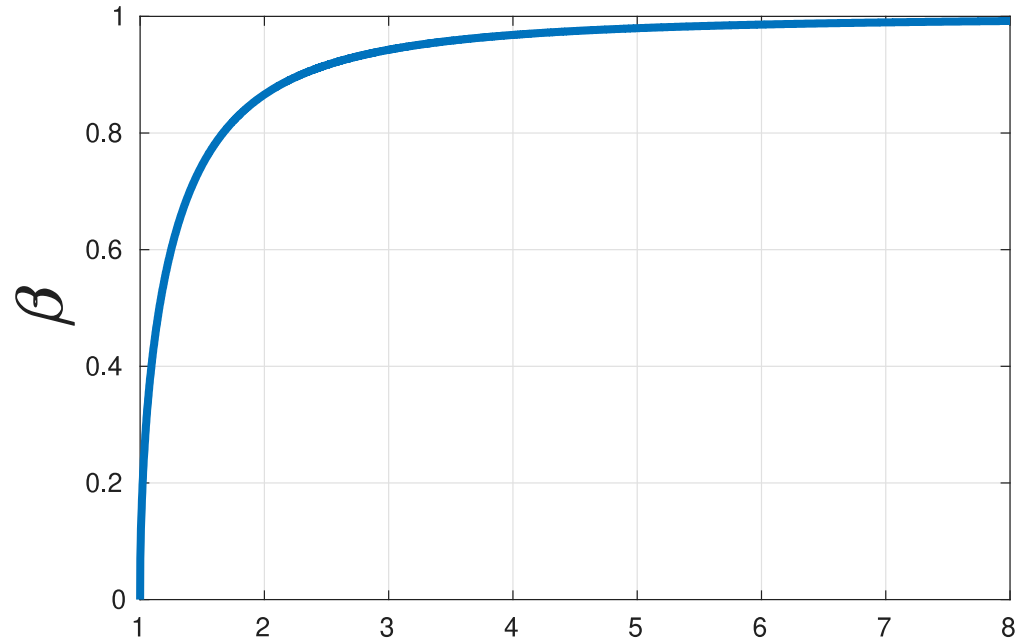
# BRIEF REVIEW OF RELATIVITY

$\beta$  is always less than 1

Velocity is almost  $c$  already for  $\gamma \sim 5$

At ESRF  $\gamma = 11800$

Increasing the electron energy from few MeV to 6 GeV does not change much the velocity.



However, there are two effects as gamma gets large:

- length contraction
- time dilation

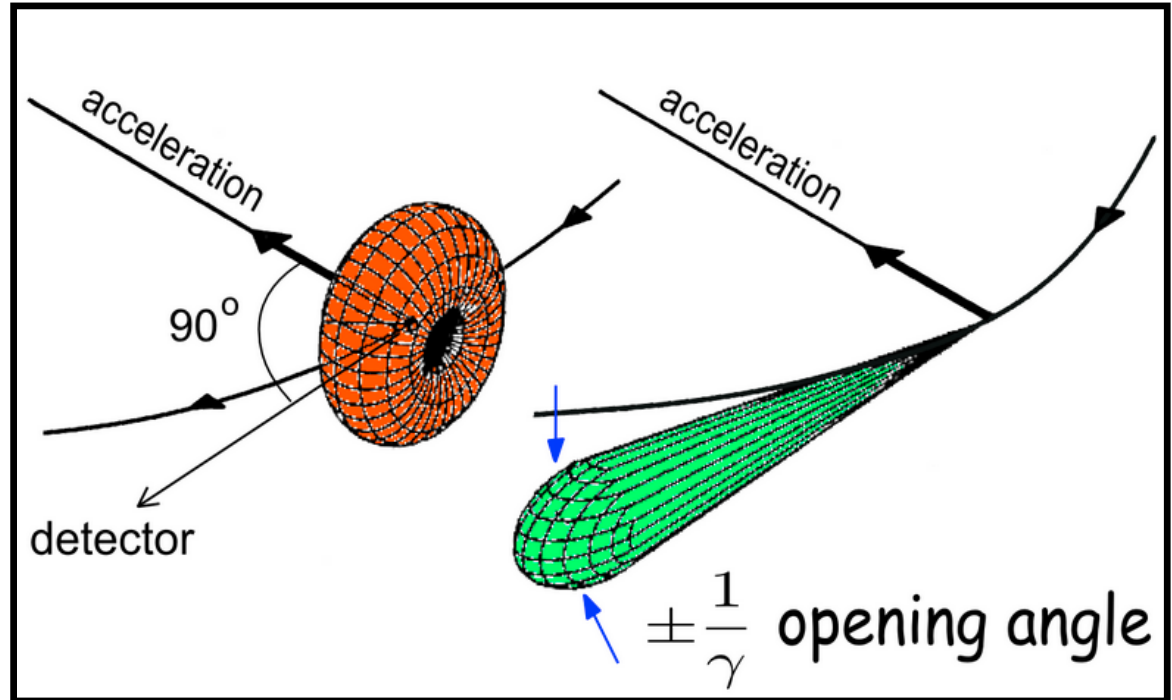
$$L = L_0 \sqrt{1 - \frac{v^2}{c^2}} = \frac{L_0}{\gamma}$$

$\gamma$

$$T = \frac{T_0}{\sqrt{1 - \frac{v^2}{c^2}}} = T_0 \gamma$$

# RADIATION EMITTED IN RELATIVISTIC CASE

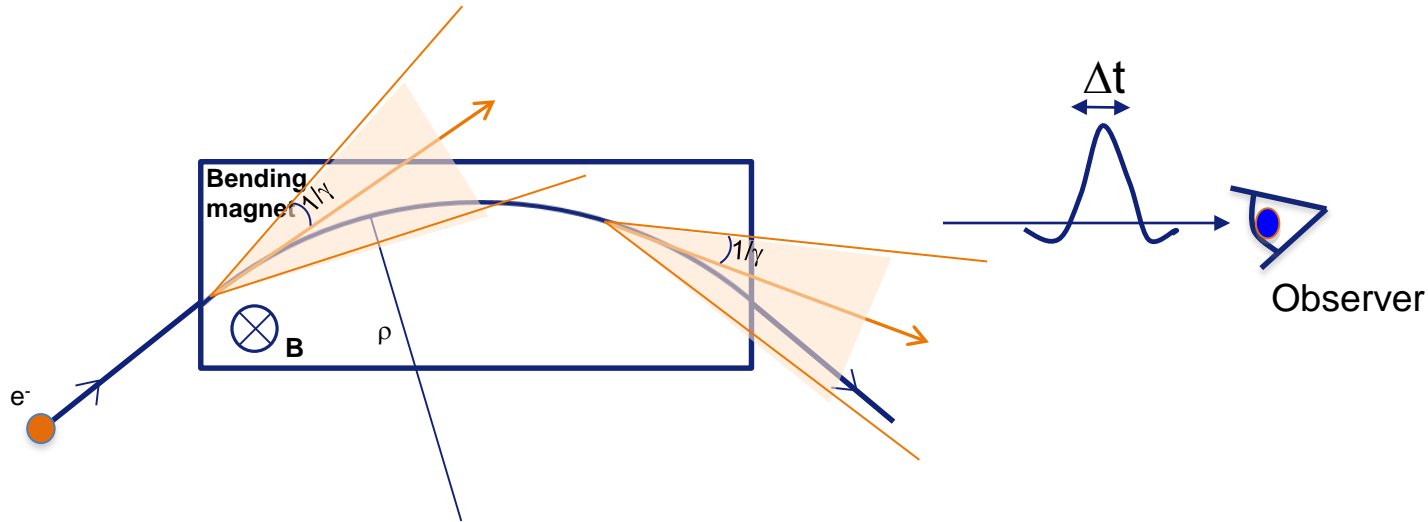
If the charged particle is moving very fast, the radiation is emitted in a cone, with aperture  $1/\gamma$



# BENDING MAGNET RADIATION: TYPICAL ENERGY

What is the **typical energy** of the radiation from a bending magnet?

First, let's see for how long an observer would see the radiation produced by an ultra-relativistic electron in a bending magnet.



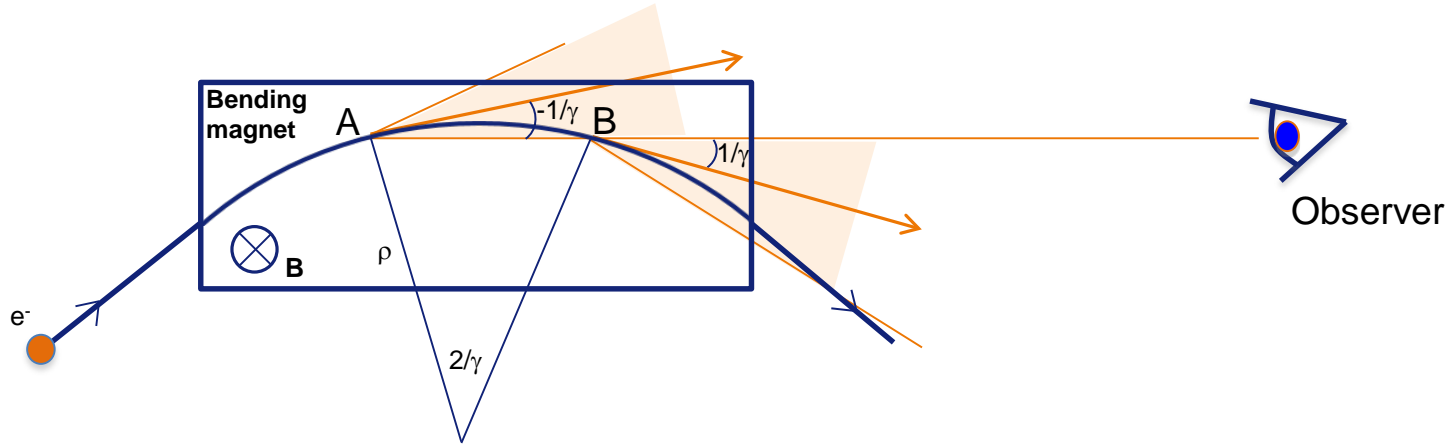
From the length  $\Delta t$  of the pulse we can estimate a typical energy  $E_{\text{typ}}$ :

$$f_{\text{typ}} = 1/\Delta t$$

$$E_{\text{typ}} = h f_{\text{typ}}$$

# BENDING MAGNET RADIATION: TYPICAL ENERGY

How long does the observer see the radiation produced by an ultra-relativistic electron in a bending magnet?



The observer will first see the radiation produced in the point A, with an angle  $-1/\gamma$ . The last photon seen will be the one produced in B with an angle  $+1/\gamma$ . The observer will therefore see radiation for an amount of time equal to the difference between the time needed by the electron to go from A to B and the time needed by the photon to go from A to B.

$$\Delta t = t_{ABe^-} - t_{AB\gamma} = \frac{\overline{AB}}{\beta c} - \frac{\overline{AB}}{c} = \frac{2\rho}{\gamma\beta c} - \frac{2\rho \sin(1/\gamma)}{c}$$

$$\sin(\theta) \simeq \theta - \frac{\theta^3}{6}, \quad \text{for } \theta \ll 1 \quad \gamma \gg 1, \quad 1 - \beta \simeq \frac{1}{2\gamma^2}$$

$$\Delta t \simeq \frac{2\rho}{\beta\gamma c} - \frac{2\rho}{c} \left( \frac{1}{\gamma} - \frac{1}{6\gamma^3} \right) = \frac{2\rho}{\beta\gamma c} \left( 1 - \beta + \frac{\beta}{6\gamma^2} \right) = \frac{4}{3} \frac{\rho}{c\gamma^3}$$

$$\omega_{typ} = \frac{2\pi}{\Delta t} = \frac{3\pi c\gamma^3}{2\rho}$$

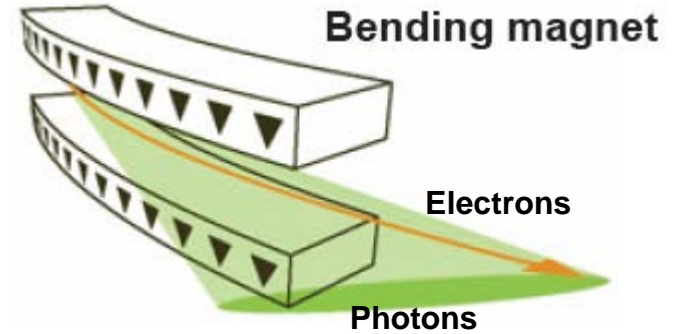
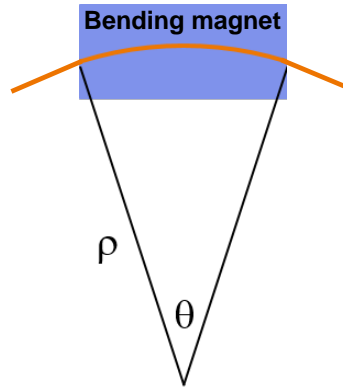
$$E_{typ} = \hbar\omega_{typ} = \frac{3\pi \hbar c}{2} \frac{\gamma^3}{\rho}$$

# BENDING MAGNET RADIATION

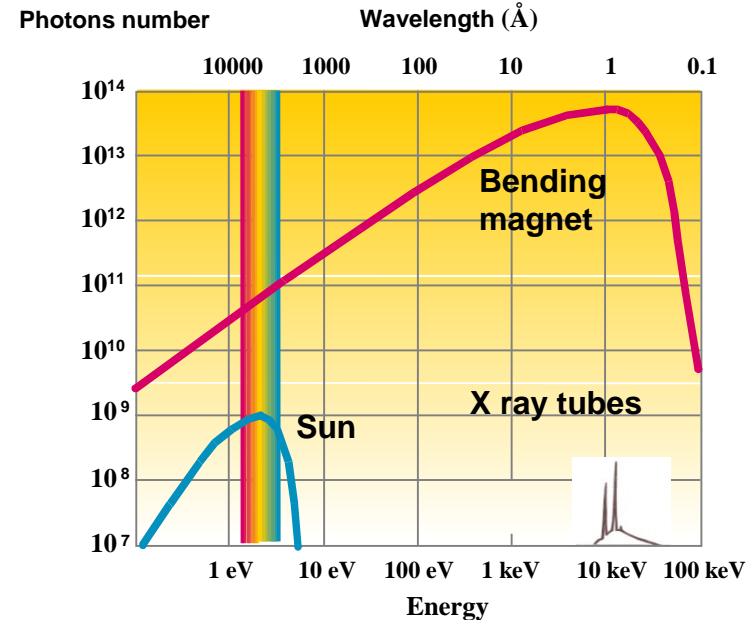
Power emitted by an ultra-relativistic particle in a bending magnet

$$P = \frac{1}{6} \frac{e^2 c}{\pi \epsilon_0} \frac{1}{\rho^2} \left( \frac{E}{mc^2} \right)^4$$

$\rho$  is the radius of curvature



## Radiation spectrum



Critical photon energy:  
half of the power is  
radiated at larger energy  
and half at lower

$$\epsilon_c = \frac{3}{2} \frac{\hbar c}{\rho} \gamma^3$$

# BENDING MAGNET RADIATION

Parameters for the EBS bending magnet beamlines

$$E = 6 \text{ GeV}$$

$$\gamma = 11800$$

$$\rho = 23.3 \text{ m}$$

$$\varepsilon_c = 21 \text{ keV}$$

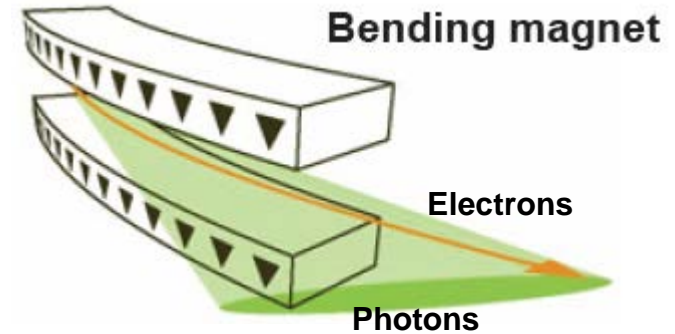
**If we put protons instead of electrons?**

$$m_p = 938 \text{ MeV}/c^2$$

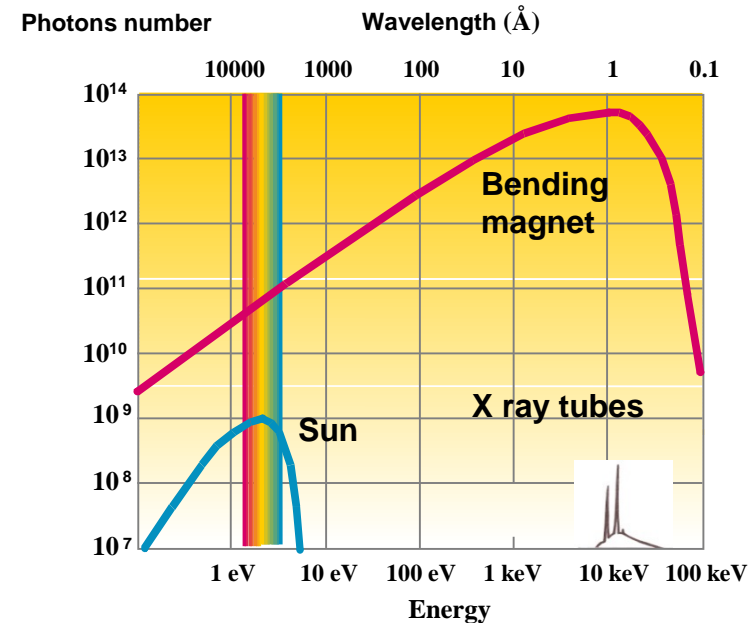
$$m_e = 0.511 \text{ MeV}/c^2$$

Radiated power would be  $10^{13}$  times smaller

$$\varepsilon_c = 3.3 \text{ } \mu\text{eV}$$

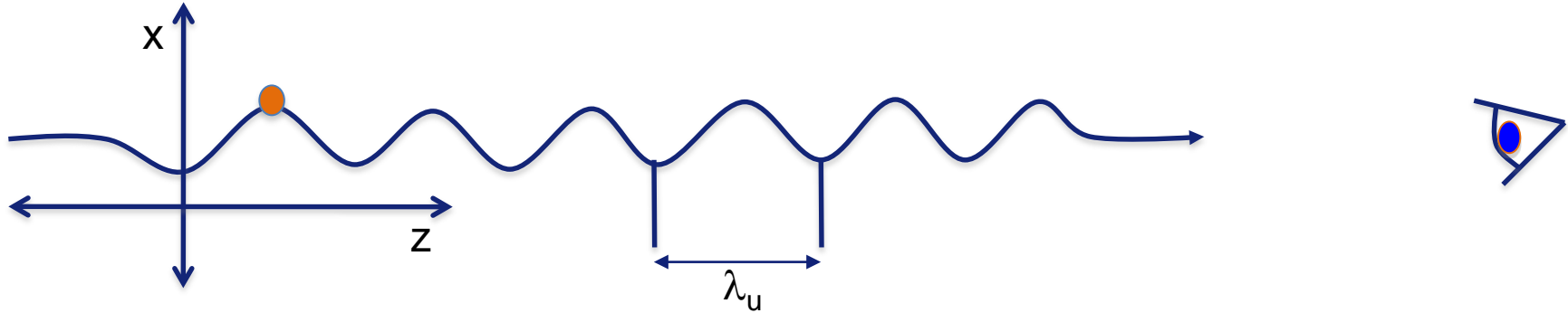


## Radiation spectrum

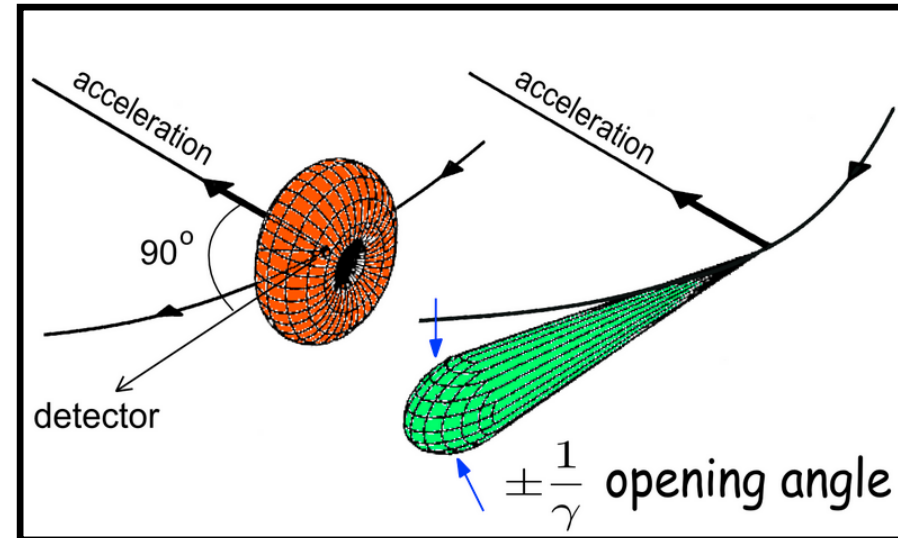


# UNDULATOR RADIATION

Consider now a single charge oscillating on x and also moving along z direction

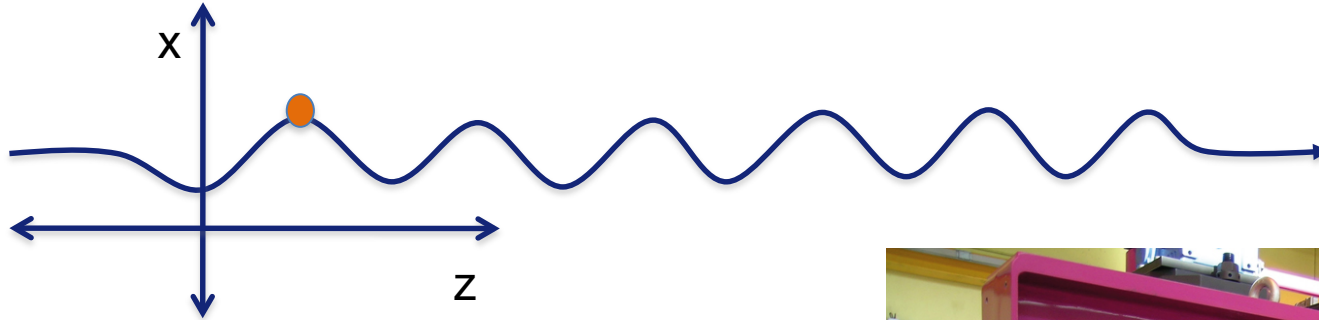


1. The wavelength is shifted by Doppler effect and, in case of relativistic speed, there is a time dilation effect.  $\lambda = \frac{\lambda_u}{\gamma^2}$
2. Pattern of radiation gets distorted from motion. For high energy gets bent into cone of angle  $1/\gamma$ .

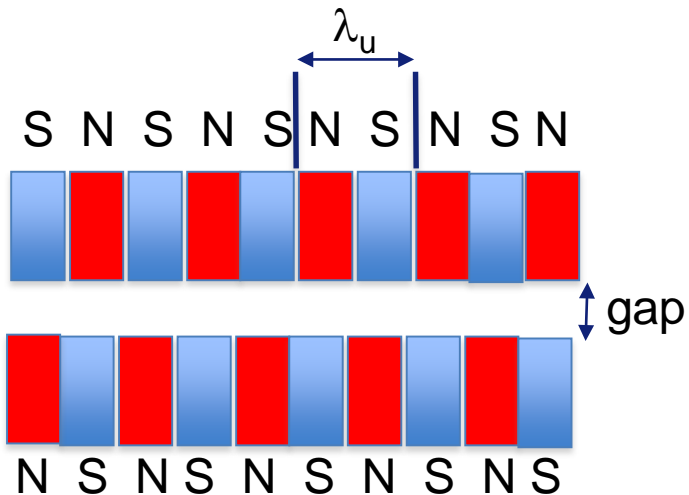


# UNDULATOR RADIATION

We can build magnets to have this kind of trajectory inside: undulators and wigglers



Undulator at ESRF





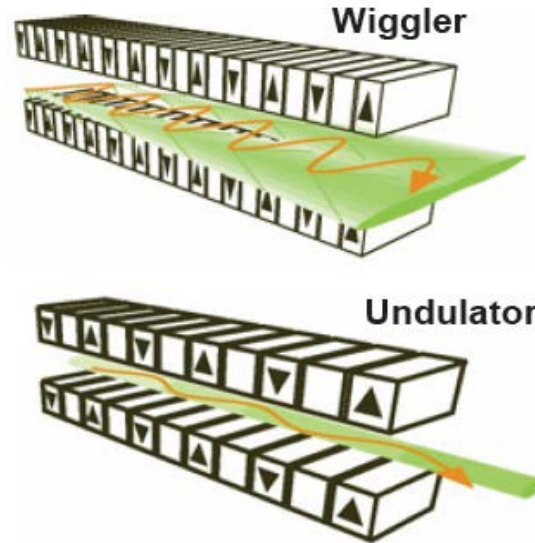
# UNDULATOR RADIATION

Undulator parameter  $K$

$$K = \frac{eB_0\lambda_u}{2\pi mc}$$

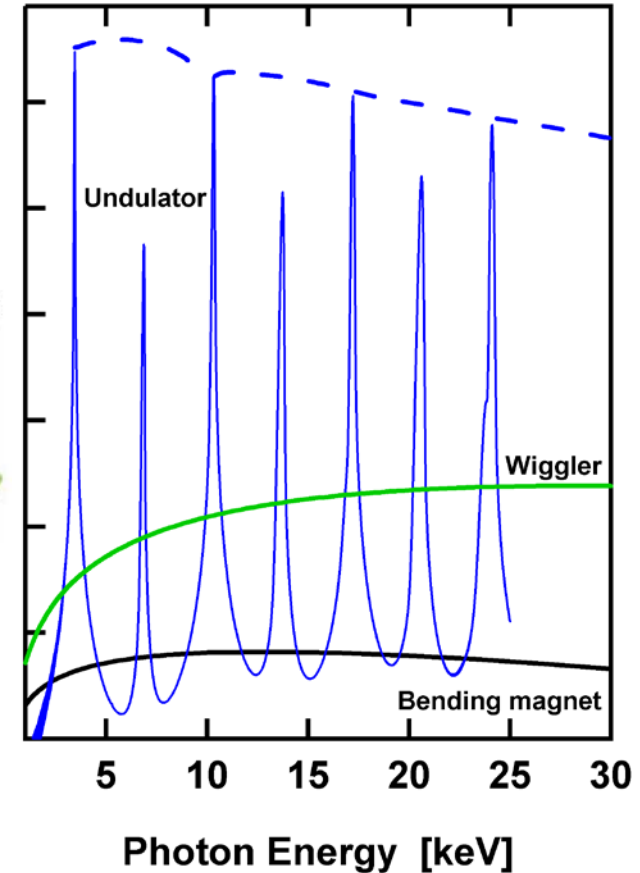
If  $K$  is large ( $K \gg 1$ ), the device is no longer called undulator, but wiggler.

$K$  can be changed by varying the gap of the undulator.

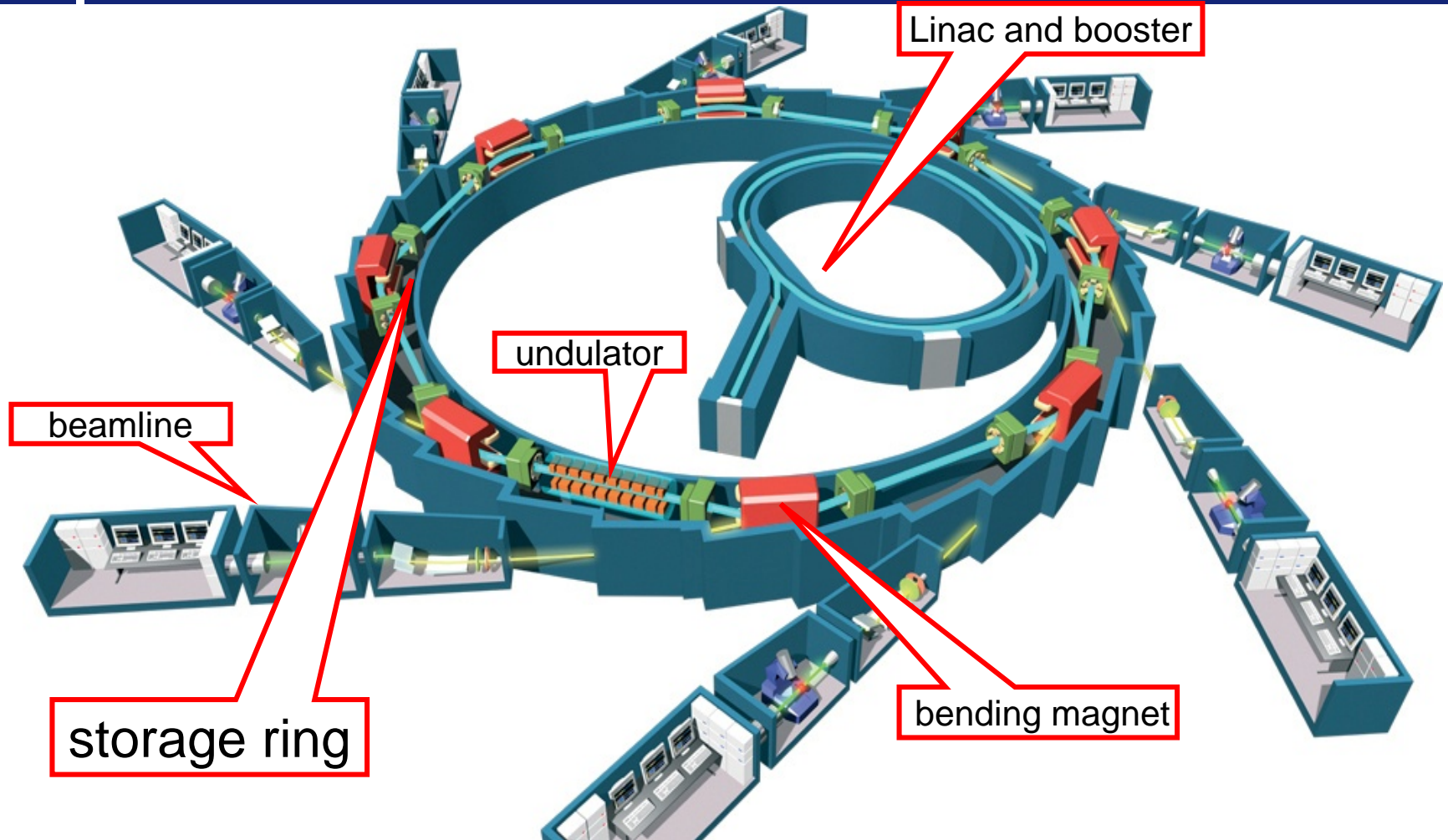


Fundamental undulator wavelength  $\lambda_1$

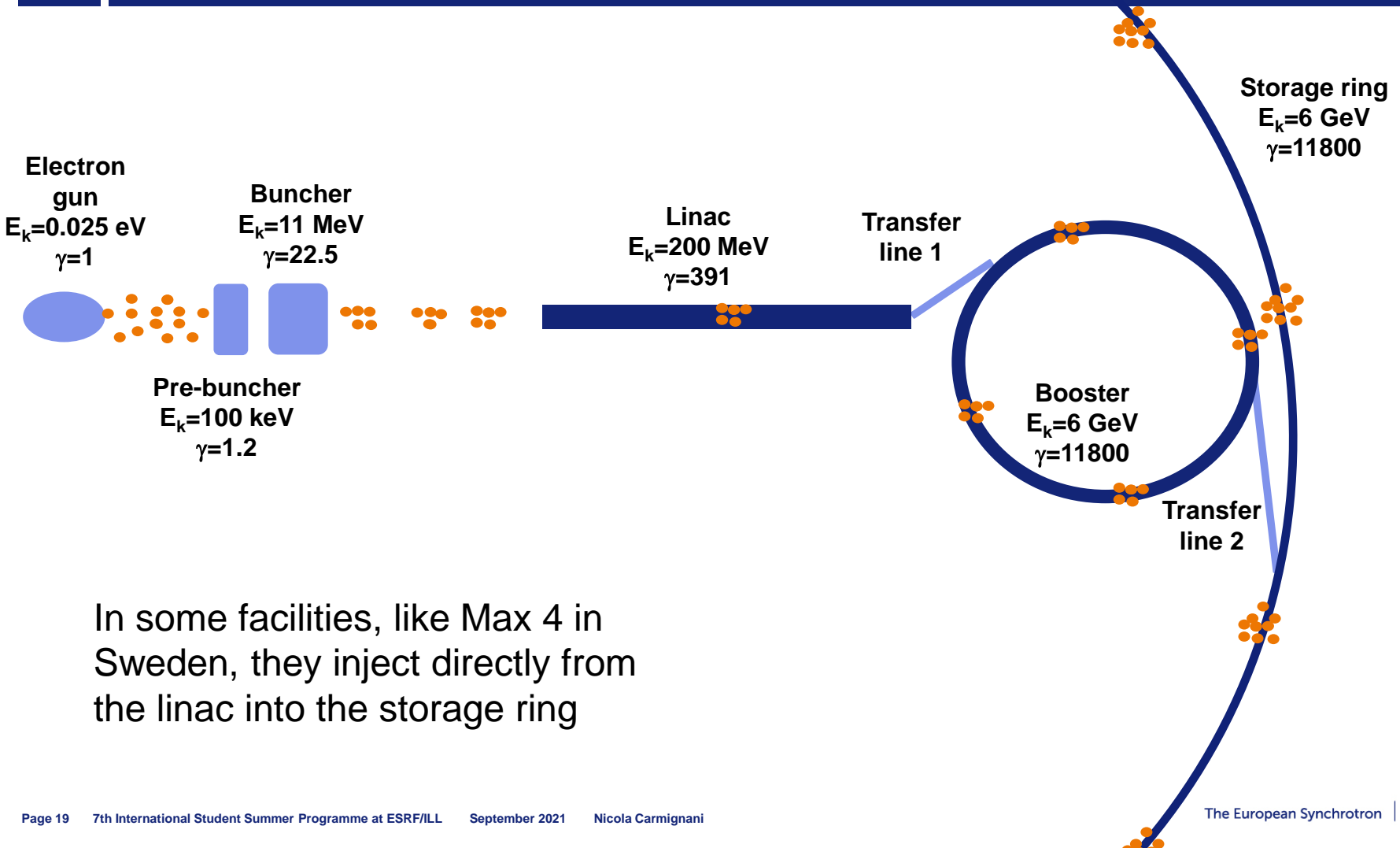
$$\lambda_1(\theta) \approx \frac{\lambda_u}{2\gamma^2} \left( 1 + \frac{K^2}{2} + (\gamma\theta)^2 \right)$$



# A TYPICAL SYNCHROTRON LIGHT SOURCE USER FACILITY



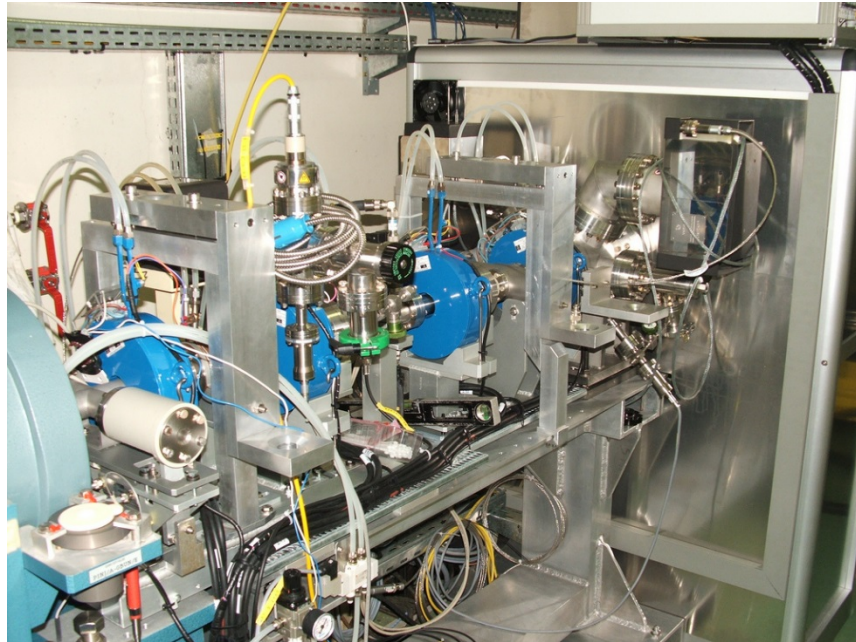
# ACCELERATION PROCESS



In some facilities, like Max 4 in Sweden, they inject directly from the linac into the storage ring

## Linac

### Electron gun, pre-buncher

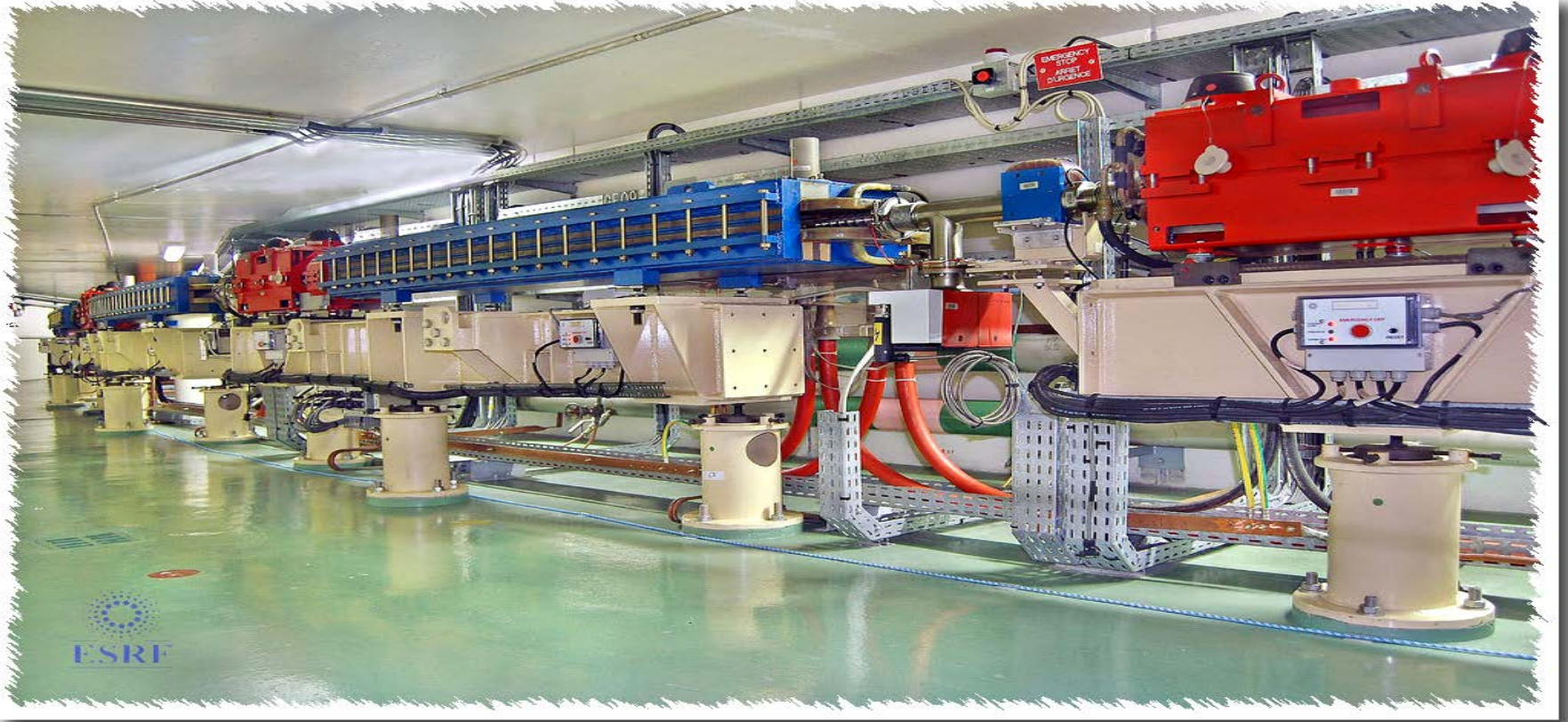


# INJECTORS: TRANSFER LINES

## Linac to booster transfer line (TL1)



## Booster



# HOW DOES A STORAGE RING WORK: MAGNETS



128 at ESRF

## **Bending magnets (dipoles)**

have uniform constant vertical magnetic field

$$B_y = B_0$$

They bend the beam and they define the circular trajectory.

At ESRF, each of the 5 pieces have a different magnetic field.

# HOW DOES A STORAGE RING WORK: MAGNETS

512 at ESRF



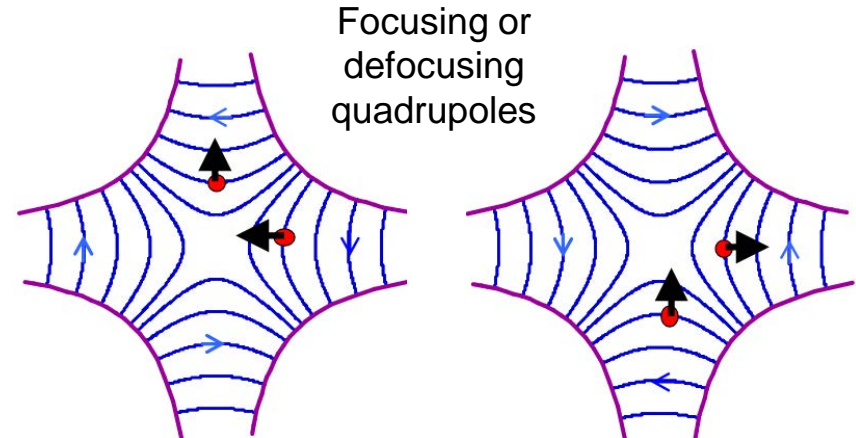
## Quadrupoles:

magnetic field is linear with the distance from the center

$$B_y = K_1 x$$

They are used to focus the beam.

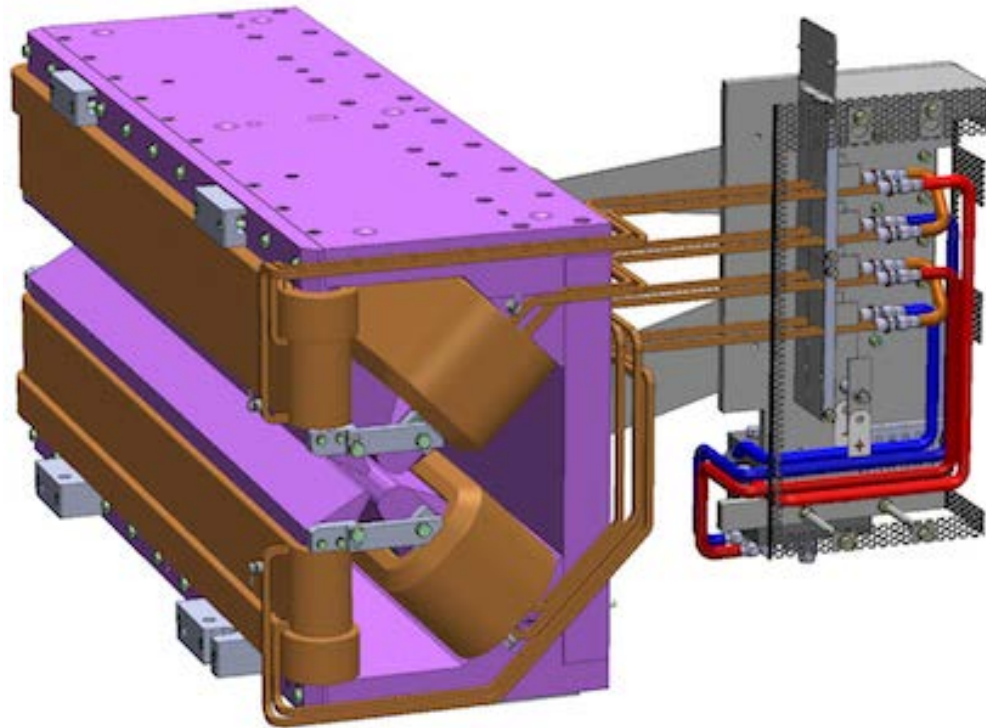
There are focusing or defocusing quadrupoles: if they focus horizontally, they defocus in vertical and vice versa (like astigmatic lenses in optics).





# HOW DOES A STORAGE RING WORK: MAGNETS

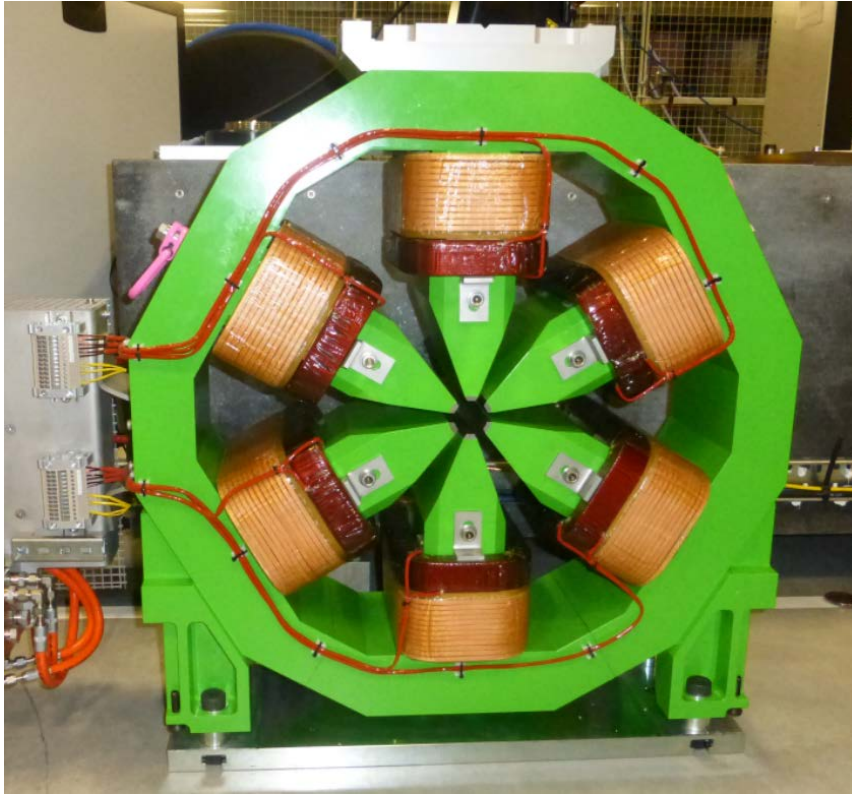
96 at ESRF



**Dipole-Quadrupoles (DQ):**  
magnetic field is the one of a dipole plus the one of a quadrupole:  
 $B_y = B_0 + K_1 x$   
They are used to bend and focus the beam.

# HOW DOES A STORAGE RING WORK: MAGNETS

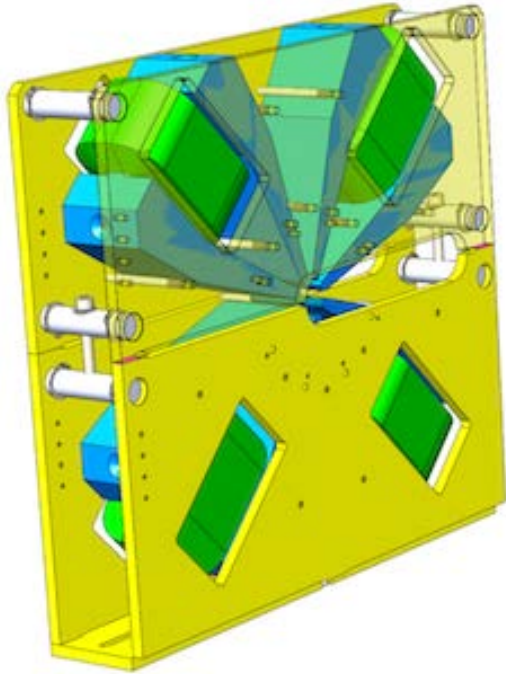
192 at ESRF



**Sextupoles:** magnetic field is quadratic with the distance from the center  
 $B_y = K_2 x^2$   
They are used to correct chromatic effects.

If correctly placed, they focus more the particles with higher energy and less the particles with lower energy. They correct chromatic effects.

# HOW DOES A STORAGE RING WORK: MAGNETS

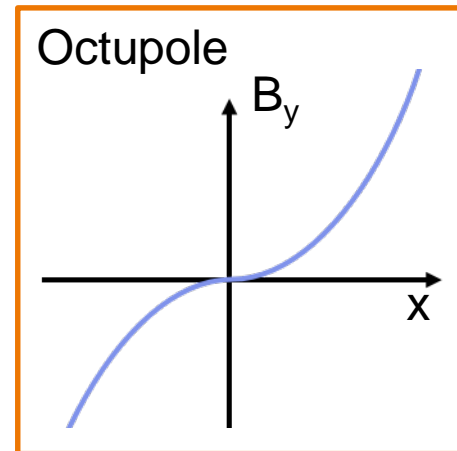
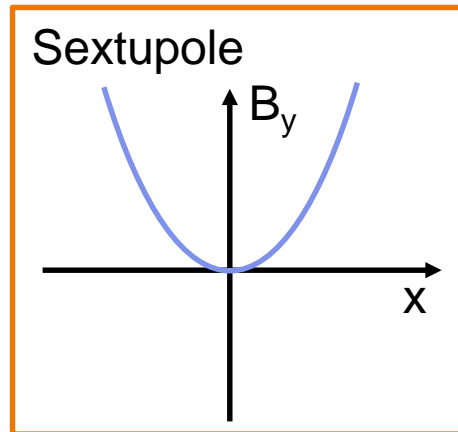
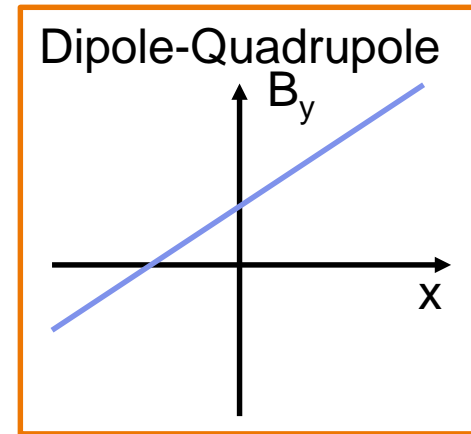
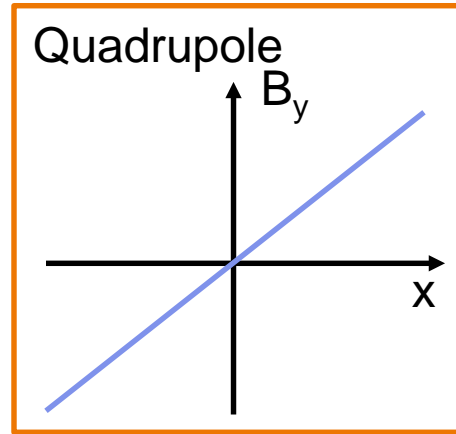
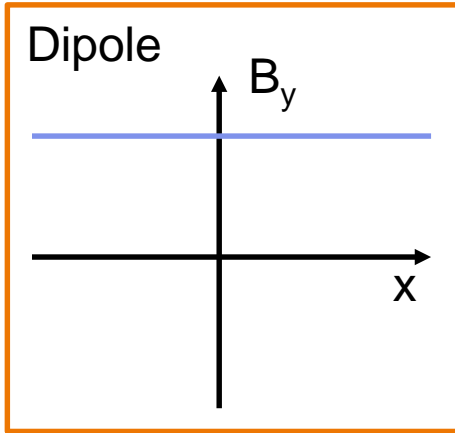


**Octupoles:**  
the field is cubic  
with distance:  
 $B_y = K_3 x^3$

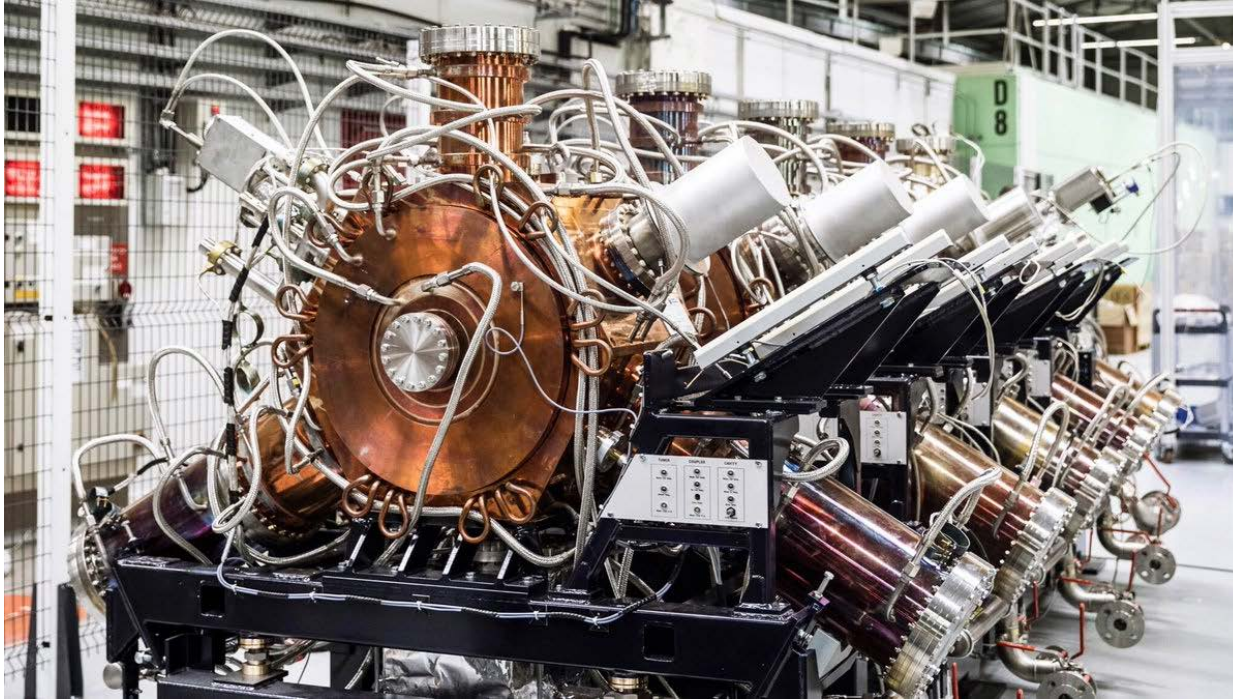


**Correctors:**  
they can provide  
horizontal dipole  
field, vertical  
dipole field,  
quadrupolar  
field.

# HOW DOES A STORAGE RING WORK: MAGNETS



# HOW DOES A STORAGE RING WORK: ACCELERATING RF CAVITIES



Most of the ESRF energy use (around 1.5 MW of power) is in these cavities

## Radio Frequency cavities

They give to the beam the energy lost in synchrotron radiation and they provide longitudinal focusing.

They produce an oscillating longitudinal electric field

$$V(t) = V_{\text{RF}} \sin(h\omega_0 t)$$

$$\omega_0 = 2\pi f_{\text{rev}}$$

At ESRF:

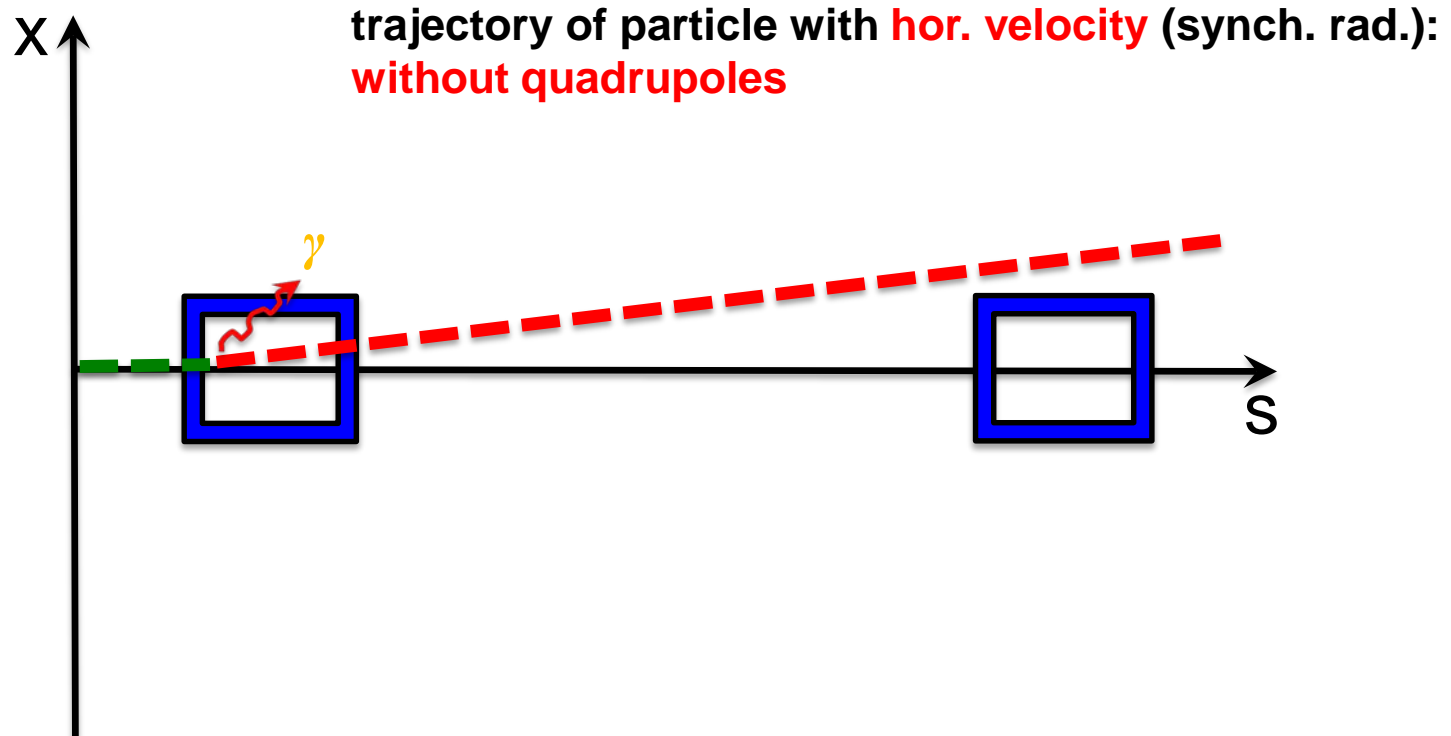
$$f_{\text{rev}} = 355 \text{ kHz}$$

$$h = 992$$

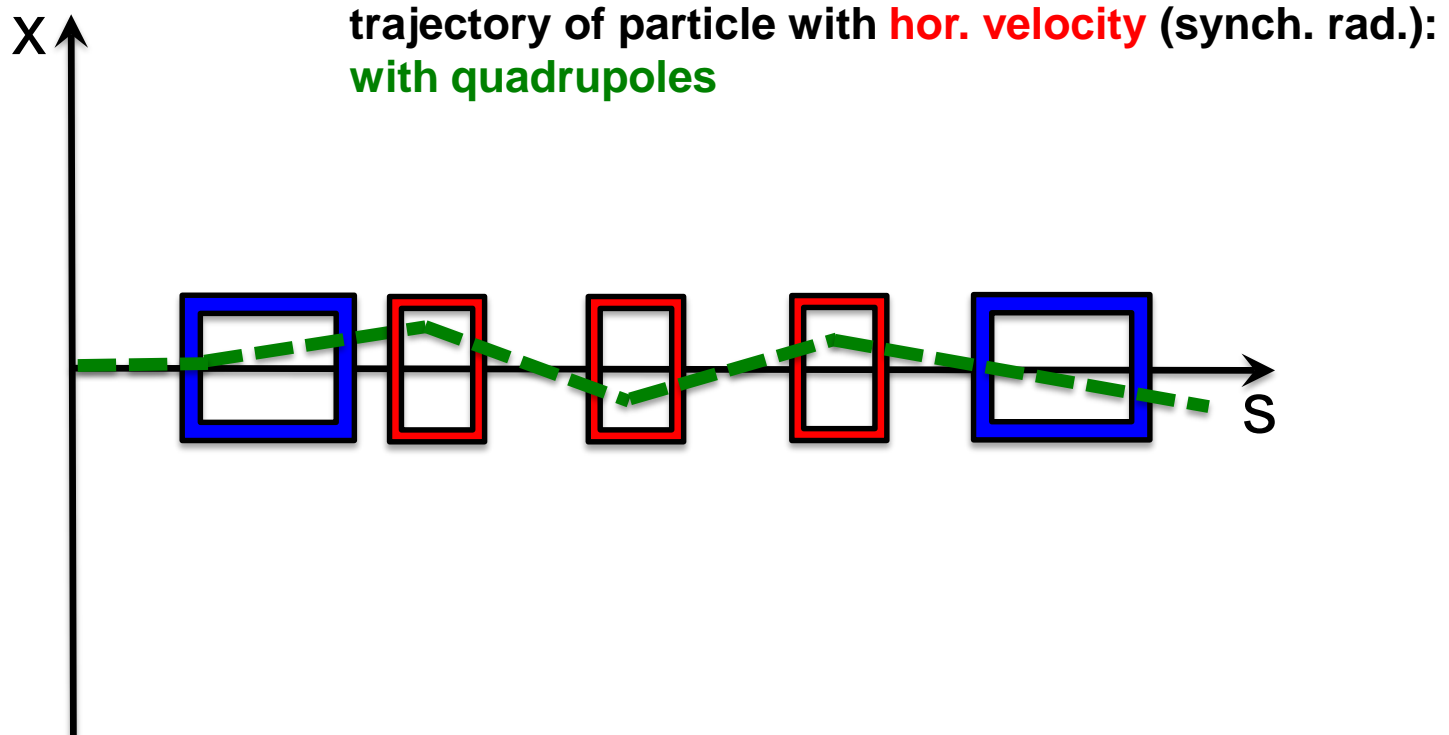
$$f_{\text{RF}} = hf_{\text{rev}} = 352 \text{ MHz}$$

$$V_{\text{RF}} = 6.5 \text{ MV}$$

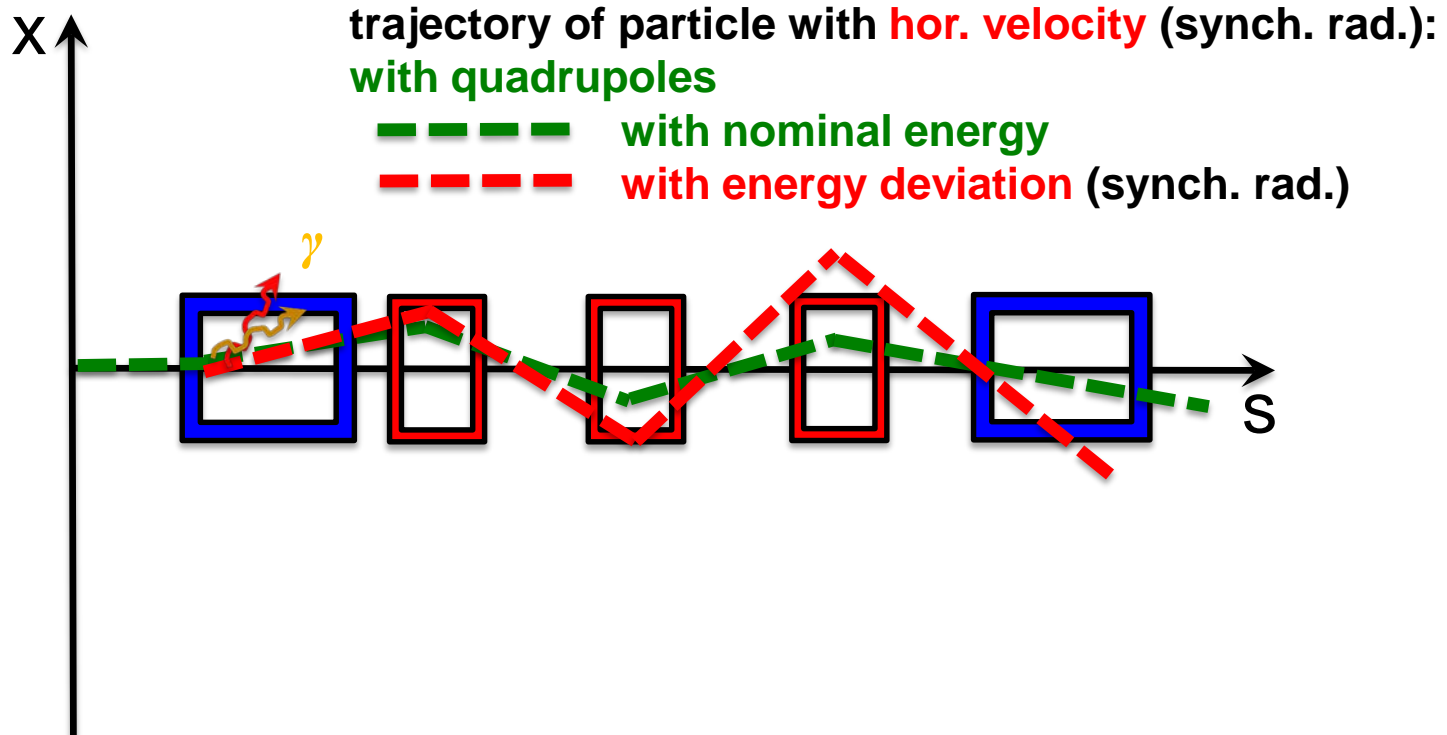
# HOW DOES A STORAGE RING WORK: DIPOLES



# HOW DOES A STORAGE RING WORK: DIPOLES, QUADRUPOLES

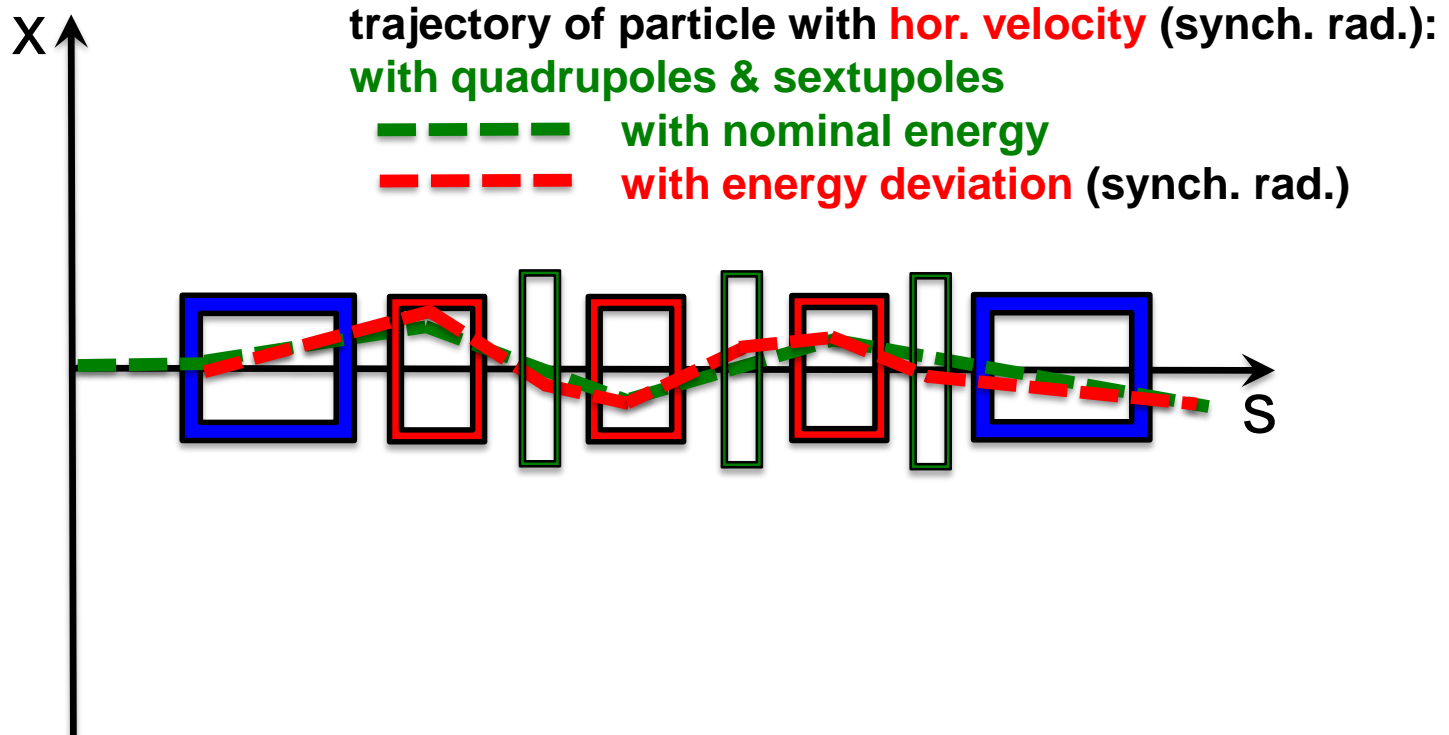


# HOW DOES A STORAGE RING WORK: DIPOLES, QUADRUPOLES

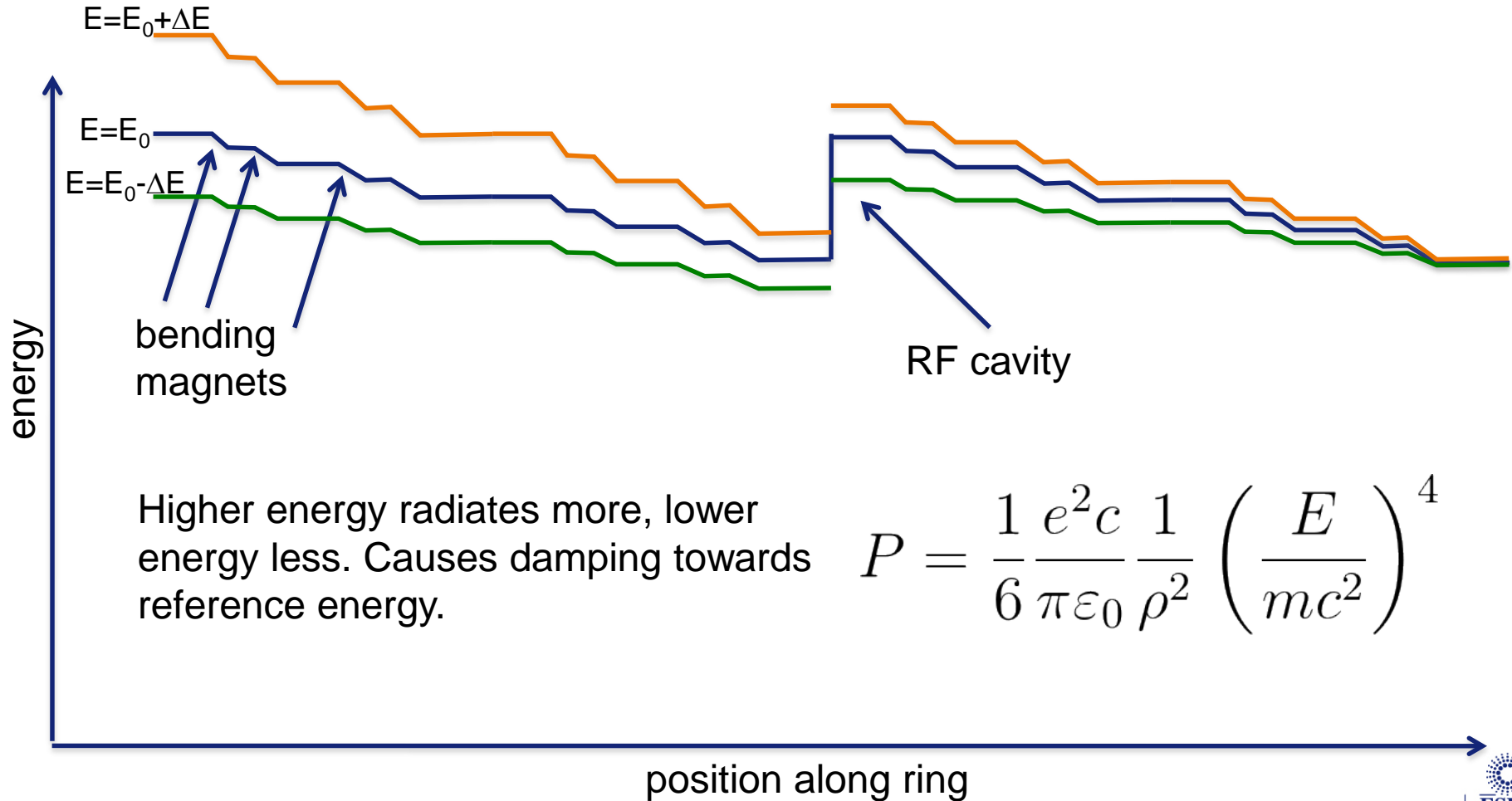




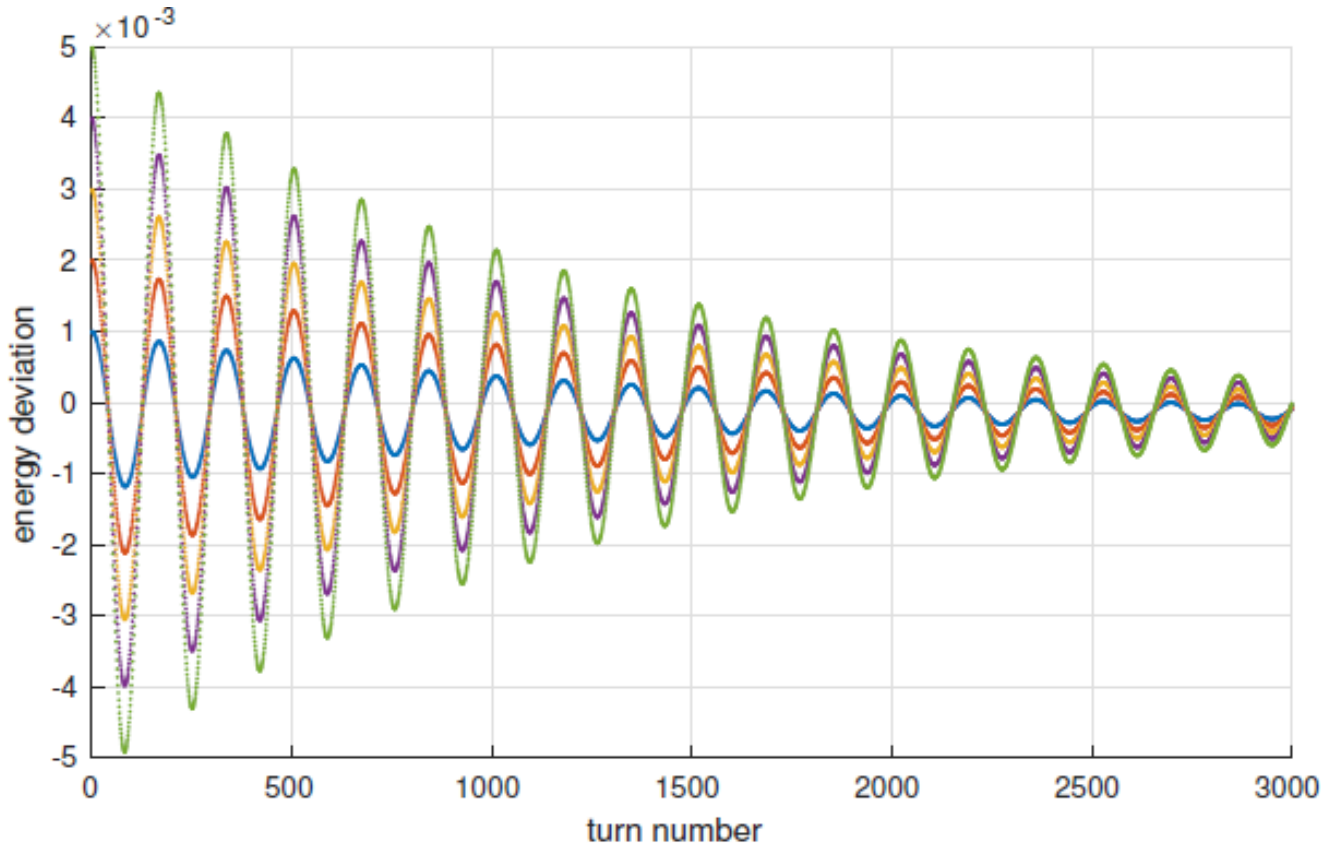
# HOW DOES A STORAGE RING WORK: DIPOLES, QUADRUPOLES, SEXTUPOLES



# DISTRIBUTION OF THE STORED ELECTRONS



# DISTRIBUTION OF THE STORED ELECTRONS



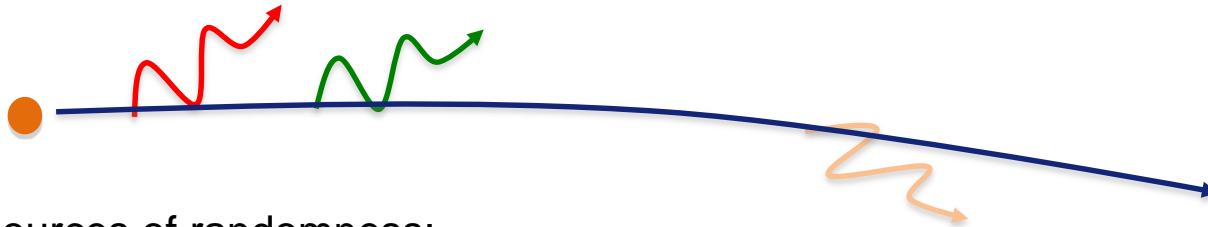
Energy oscillations  
(**synchrotron oscillations**)  
and radiation damping.

So what defines the size of  
the electron beam?

# DISTRIBUTION OF THE STORED ELECTRONS

Question: if we have radiation damping, why isn't the electron beam size zero?

Answer: because of the graininess of the photon emission



Two sources of randomness:

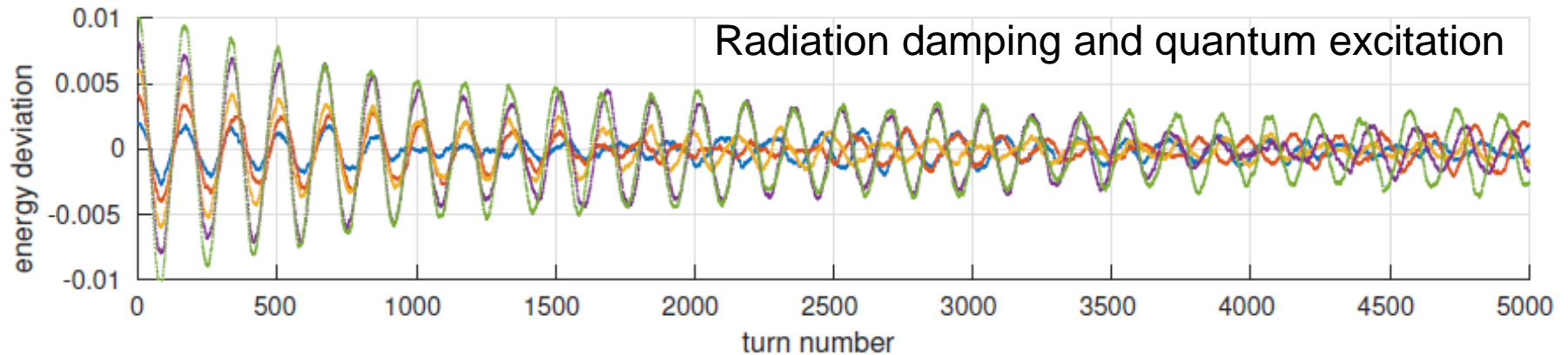
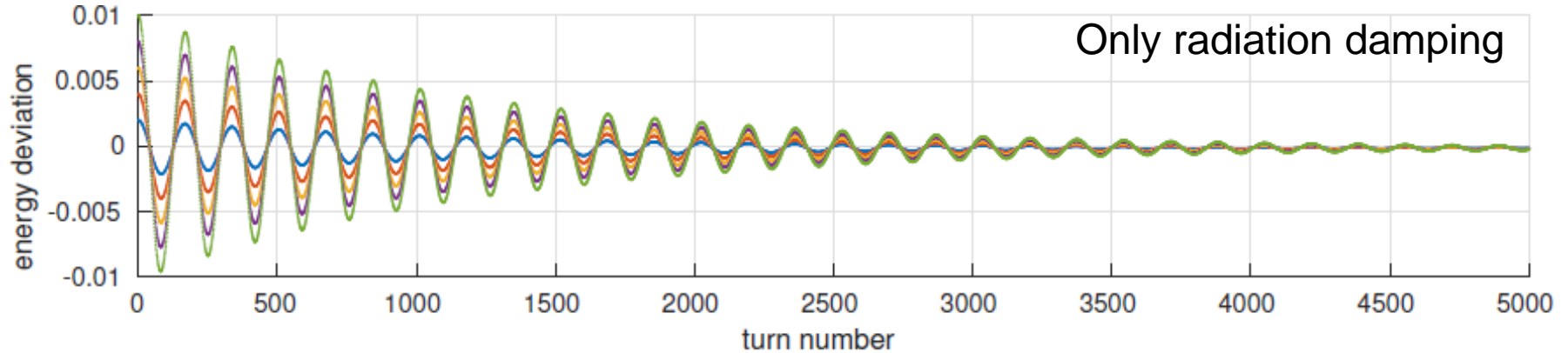
- emission time of photons: Poisson process
- energy of the photon emitted

This causes a diffusion process in energy, that is converted to a transverse diffusion by the dipoles (higher energy electrons are bended less than lower energy ones).

For ESRF, each electron emits only a few hundreds photons per turn.

Energy loss per turn  $U_0 = 2.53$  MeV

# DISTRIBUTION OF THE STORED ELECTRONS



# Results of damping and diffusion

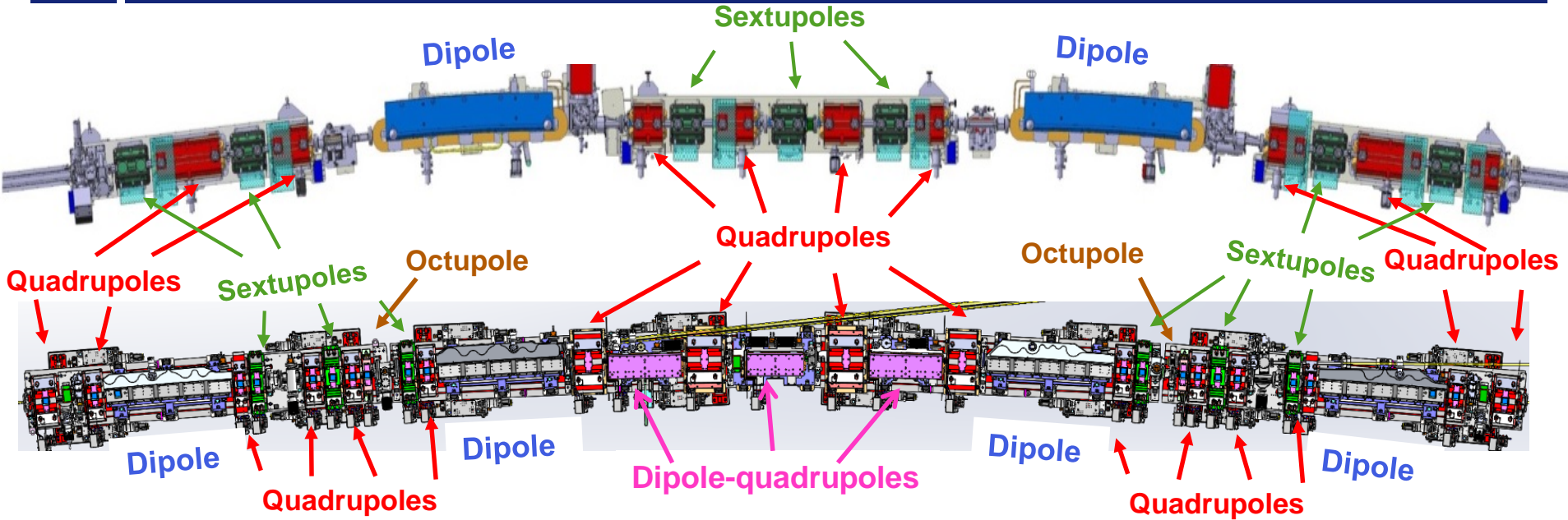
The electron beam reaches a unique Gaussian distribution, independent of how one injects into the ring.

This is a major difference between electron synchrotrons and proton synchrotrons (e.g. LHC)

By careful choice of where the dipoles and quadrupoles are, one can reduce the size of this equilibrium beam size (emittance = beam size in phase space).  
So called “low emittance ring design”.

In fact, due to developments in lattice design, ESRF has completely replaced the storage ring in 2019 to reduce the electron beam emittance. 4nm rad -> 130 pm rad

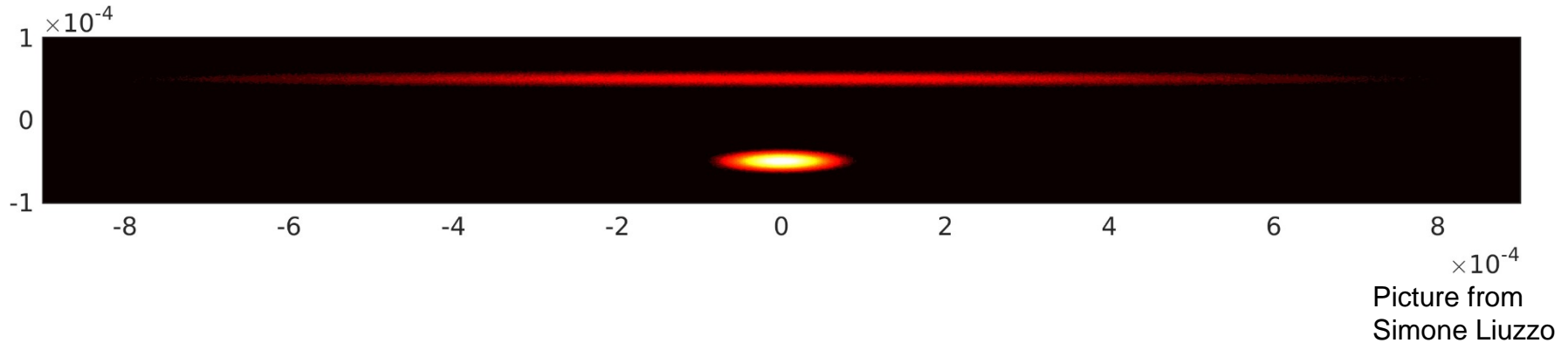
# THE EBS UPGRADE AT ESRF



The ESRF **Extremely Brilliant Source** (EBS) has 31 magnets per cell instead of 17 of the old machine. Free space between magnets (total for one cell): **3.4 m** instead of **8 m**!

Magnet type	ESRF	ESRF EBS
Dipoles	64	128
Dipole-Quadrupoles	0	96
Quadrupoles	256	512
Sextupoles	224	192
Octupoles	0	64
Correctors	0	96

## x-y beam distribution for old and new machine.



Electron beam horizontal emittance has been reduced by a factor 30 and the brightness is increases by the same factor.

$$\begin{aligned}\varepsilon_{\text{ESRF}} &= 4 \text{ nm rad} \\ \varepsilon_{\text{EBS}} &= 130 \text{ pm rad}\end{aligned}$$



## **Electromagnetism:**

J. D. Jackson, “Classical Electrodynamics”, Wiley.

R. P. Feynman, “The Feynman Lectures on Physics, Volume 2”, [www.feynmanlectures.caltech.edu](http://www.feynmanlectures.caltech.edu).

## **Synchrotron radiation:**

A. Hofmann, “Characteristics of Synchrotron Radiation”, proceedings of the CERN accelerator school 1998.

P. Elleaume, “Undulators and wigglers”, Taylor & Francis.

## **Particle accelerator physics:**

H. Wiedemann, “Particle accelerator physics”, Springer.

A. Chao, “Handbook of Accelerator Physics and Engineering”, World scientific.

S. Y. Lee, “Accelerator physics”, World scientific.

M. Sands, “The physics of electron storage rings, an introduction”, SLAC report no. 121.

## **ESRF EBS upgrade:**

The orange book, ESRF upgrade programme phase I (2015-2022) Technical Design Study.

MANY THANKS FOR YOUR ATTENTION

